

Operator-Theoretic Approaches for Coherent Feature Extraction in Complex Systems

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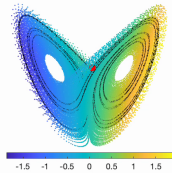
Institute for Mathematical and Statistical Innovation
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Collaborators: S. Das, G. Froyland, B. Lintner, M. Pike, J. Slawinska

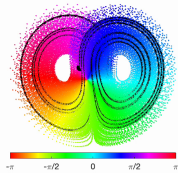


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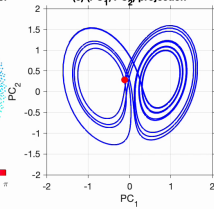
(a) PC1 on L63 attractor



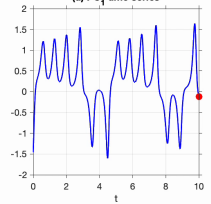
(b) (PC_1, PC_2) angle on L63 attractor



(c) (PC_1, PC_2) projection

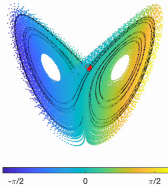


(d) PC1 time series

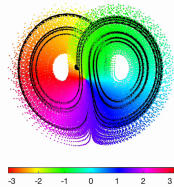


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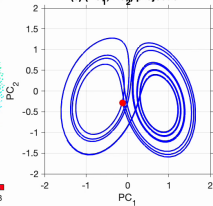
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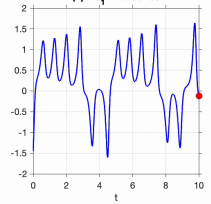
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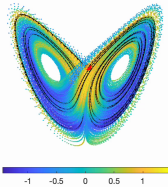
(c) (PC_1, PC_2) projection



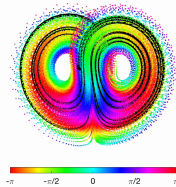
(d) PC_1 time series



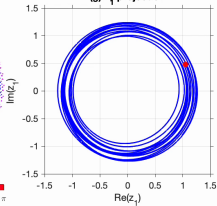
(e) $\text{Re}(z_1)$ on L63 attractor



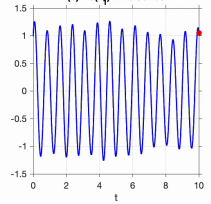
(f) z_1 angle on L63 attractor



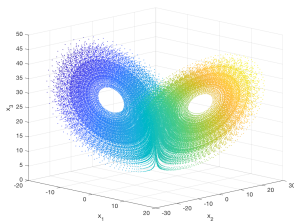
(g) z_1 projection



(h) $\text{Re}(z_1)$ time series



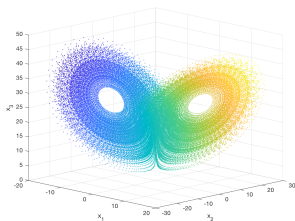
Operator-theoretic description of dynamical systems



A dynamical flow $\Phi^t : X \rightarrow X$ acts on **spaces of observables** through intrinsically linear operators (**Koopman operators**):

$$\mathcal{F} = \{f : X \rightarrow \mathbb{C}\},$$
$$U^t : \mathcal{F} \rightarrow \mathcal{F}, \quad U^t f(x) = f(\Phi^t(x)).$$

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$$f = g + h \implies U^t f(x) = g(\Phi^t(x)) + h(\Phi^t(x)) = U^t g(x) + U^t h(x),$$
$$f = cg \implies U^t f(x) = f(\Phi^t(x)) = cg(\Phi^t(x)) = cU^t g(x).$$

Objectives

- 1 Identify **coherent observables** under the dynamics (*spectral approximation of operators*).
 - 2 Perform **forecasting** of observables (*pointwise approximation of operators*).
- Methods should be **data-driven**, i.e., only utilize information from a time-ordered sequence of measurements.
 - Challenges due to infinite dimensionality include the presence of **unbounded operators** and **continuous spectra**.

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 - Challenges due to infinite dimensionality include the presence of **unbounded operators** and **continuous spectra**.
 - Operator-theoretic formalism originally introduced in the 1930s (Koopman 1931; von Neumann 1932).
 - More recently, there has been significant interest in data-driven approaches (Dellnitz & Junge 1999; Dellnitz & Froyland 2000; Mezić & Banaszuk 2004; Mezić 2005; Rowley et al. 2009; Schmidt 2010; Froyland et al. 2014; Berry et al. 2015; Froyland 2015; G. et al. 2015; Williams et al. 2015; Klus et al. 2016; Brunton et al. 2017; Wu & Noé 2017; Korda et al. 2018; Das et al. 2019, 2020, 2021; G. 2019, 2021; Thiede et al. 2019; Alexander & G. 2020; Berry et al. 2020...).

Coherent observables through Koopman spectral analysis

A nonzero observable $z : X \rightarrow \mathbb{C}$ is said to be a **Koopman eigenfunction** if there exists $\lambda_t \in \mathbb{C}$ such that

$$U^t z = \lambda_t z.$$

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- Restrict attention to eigenfunctions $z \in L^2(\mu)$, where μ is an **invariant measure** of the dynamics.
- In that case, $\lambda_t = e^{i\omega t}$, where $\omega \in \mathbb{R}$ is an **eigenfrequency** of the dynamical system.
- Koopman eigenfunctions in measure-preserving systems form a distinguished class of observables evolving periodically under the dynamics at intrinsic frequencies, even if the flow map Φ^t is aperiodic:

$$U^t z = e^{i\omega t} z.$$

Example: Rectification of variable-frequency oscillator

$$\begin{aligned}\dot{\theta} &= v(\theta) \equiv \alpha(1 + \beta \sin \theta), \\ \theta &\in S^1, \quad \alpha > 0, \quad 0 \leq \beta < 1\end{aligned}$$

- Koopman eigenfunctions:

$$\begin{aligned}z_j(\theta) &= e^{i\omega_j h(\theta)}, \quad \omega_j = \frac{2\pi j}{T}, \quad j \in \mathbb{Z}, \\ h(\theta) &= \int_0^\theta \frac{1}{v(\vartheta)} d\vartheta, \quad T = h(2\pi).\end{aligned}$$

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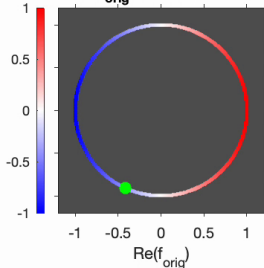
- Viewed as features $z_j : S^1 \rightarrow \mathbb{C}$, the Koopman eigenfunctions **rectify** the oscillation:

$$\begin{array}{ccc} S^1 & \xrightarrow{\Phi^t} & S^1 \\ z_j \downarrow & & \downarrow z_j \\ \mathbb{C} & \xrightarrow{e^{i\omega_j t}} & \mathbb{C} \end{array}$$

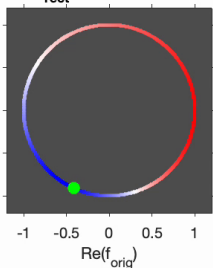
Example: Rectification of variable-frequency oscillator

$t = 14.14$

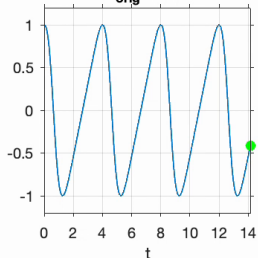
(a) $\text{Re}(f_{\text{orig}})$ on original space



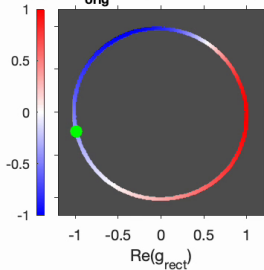
(b) $\text{Re}(g_{\text{rect}} \circ h)$ on original space



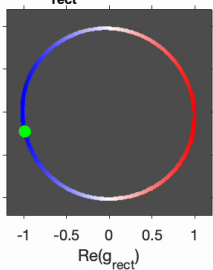
(c) $\text{Re}(f_{\text{orig}})$ time series



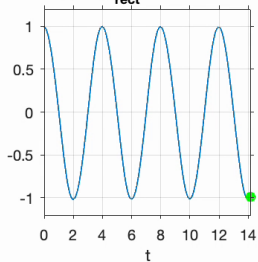
(d) $\text{Re}(f_{\text{orig}} \circ h^{-1})$ on rectified space



(e) $\text{Re}(g_{\text{rect}})$ on rectified space



(f) $\text{Re}(g_{\text{rect}})$ time series



Coherent observables through Koopman spectral analysis

An observable $z \in L^2(\mu)$ is said to be an ϵ -approximate **Koopman eigenfunction** if there exists $\nu_t \in \mathbb{C}$ such that

$$\|U^t z - \nu_t z\| < \epsilon \|z\|.$$

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- A Koopman eigenfunction is an ϵ -approximate eigenfunction for every $\epsilon > 0$.
- A natural notion of **coherence** is to require that z is an ϵ -approximate eigenfunction for “small” ϵ , and t lying in a “large” time interval.

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Given: Time series data y_0, \dots, y_{N-1} , with $y_n = F(x_n)$, of an observable $F : X \rightarrow \mathbb{R}^d$, sampled along a dynamical trajectory, $x_n = \Phi^{n\Delta t}(x_0)$.

- 1 Map the data into **delay-coordinate space**,

$$y_n \mapsto \tilde{y}_n = (y_n, y_{n-1}, \dots, y_{n-Q+1}) \in \mathbb{R}^{dQ}.$$

- 2 Introduce a **kernel function** $k_Q : \mathbb{R}^{dQ} \times \mathbb{R}^{dQ} \rightarrow \mathbb{R}$, and compute the associated $N \times N$ kernel matrix,

$$\mathbf{K}_Q = [k_Q(\tilde{y}_m, \tilde{y}_n)], \quad 0 \leq m, n \leq N-1.$$

- 3 Compute eigenvalues and eigenvectors of \mathbf{K}_Q ,

$$\mathbf{K}_Q \vec{\phi}_j = \lambda_j \vec{\phi}_j.$$

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Result [Das & G. 19; G. 21]: *For a suitable choice of kernel k_Q the following hold:*

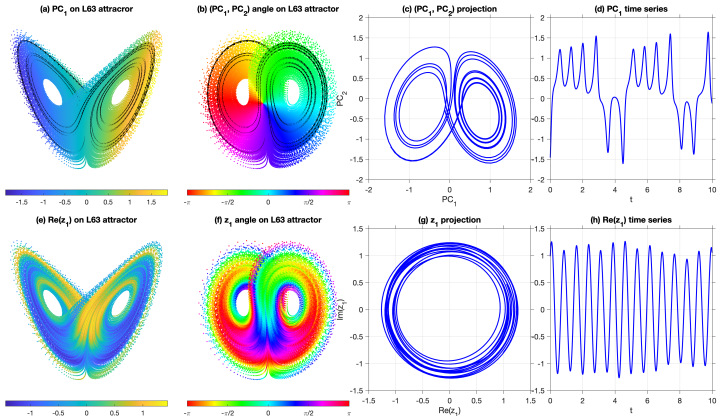
- 1 As $N \rightarrow \infty$, $\vec{\phi}_j \in \mathbb{R}^N$ converges to a continuous eigenfunction $\phi_j : X \rightarrow \mathbb{R}$ of a kernel integral operator K_Q on $L^2(\mu)$.
- 2 If the eigenvalue corresponding to ϕ_j is sufficiently isolated in the spectrum of K_Q , then the complex-valued observable $z = \phi_j + i\phi_{j+1}$ is an ϵ_t -approximate eigenfunction of U^t for a bound

$$\epsilon_t \sim \frac{t}{\lambda_j Q}.$$

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Galerkin approximation of the generator (G. et al. 17; G. 19; Thiede et al. 19)

$$V : D(V) \rightarrow L^2(\mu), \quad D(V) \subset L^2(\mu)$$
$$Vf = \lim_{t \rightarrow 0} \frac{U^t f - f}{t}, \quad V^* = -V, \quad U^t = e^{tV}$$

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- The matrix elements of the generator can be approximated in the kernel eigenfunction basis using **temporal finite differences**:

$$\begin{aligned} \langle \phi_i, V \phi_j \rangle &= \int_{\mathcal{X}} \phi_i(x) V \phi_j(x) d\mu(x) \\ &= \lim_{\Delta t \rightarrow 0} \lim_{N \rightarrow \infty} \frac{\vec{\phi}_i \cdot \mathbf{U}^{\Delta t} \vec{\phi}_j - \delta_{ij}}{N \Delta t} \\ &\approx \frac{1}{N \Delta t} \sum_{n=0}^{N-1} \phi_{n,i} (\phi_{j,n+1} - \phi_{j,n}). \end{aligned}$$

Galerkin approximation of the generator (G. et al. 17; G. 19; Thiede et al. 19)

$$V_\epsilon z_j = \gamma_j z_j$$
$$V_\epsilon = V - \epsilon \Delta, \quad \epsilon > 0$$

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- Δ is a **diffusion operator**, constructed from the eigenvalues/eigenfunctions of the kernel integral operator:

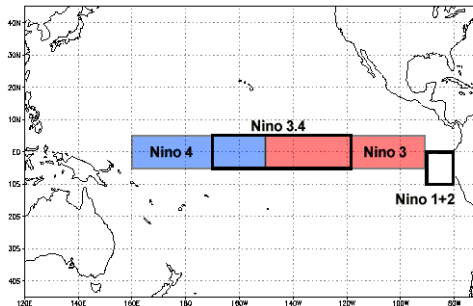
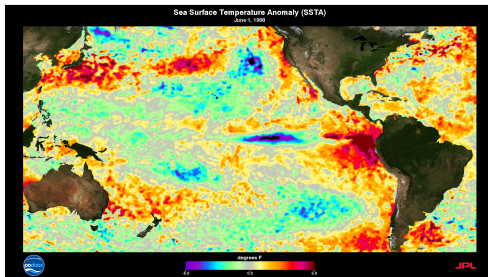
$$\Delta \phi_i = \eta_i \phi_i, \quad \eta_i = \frac{1}{\lambda_i} - 1, \quad i \in \{0, 1, \dots, \}.$$

- The eigenvalue problem for the regularized generator V_ϵ is approximated by a matrix eigenvalue problem,

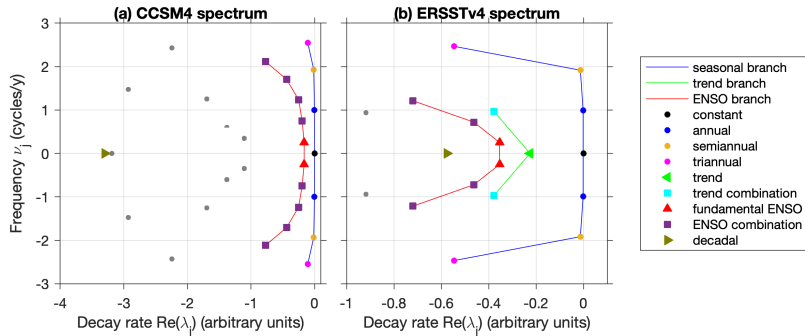
$$(\mathbf{V} - \epsilon \mathbf{\Delta}) \vec{z}_j = \gamma_j \mathbf{G} \vec{z}_j.$$

- $\text{Im } \gamma_j$: Approximate Koopman eigenfrequency.
- $\text{Re } \gamma_j$: Dirichlet energy (decay rate) of z_j .

El Niño Southern Oscillation (ENSO)

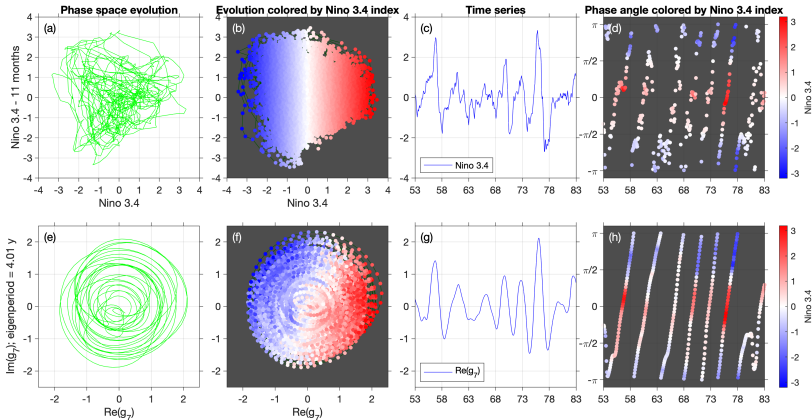


Koopman spectra in models and observations (Froyland et al. 21)



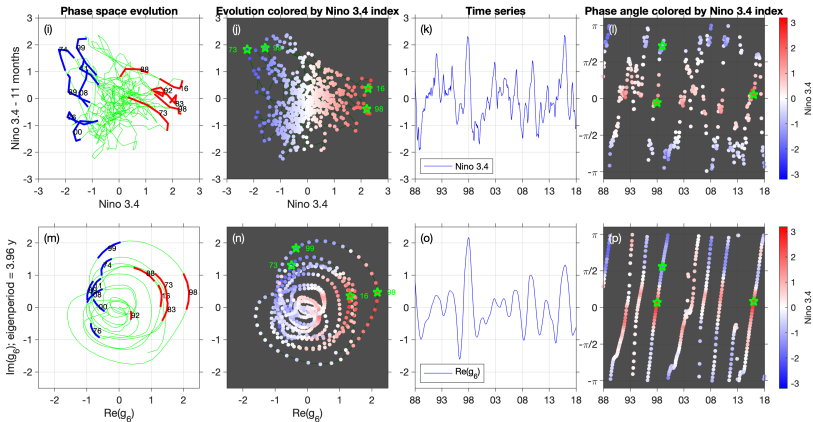
ENSO lifecycle recovered from approximate Koopman eigenfunctions

CCSM4 ENSO lifecycle

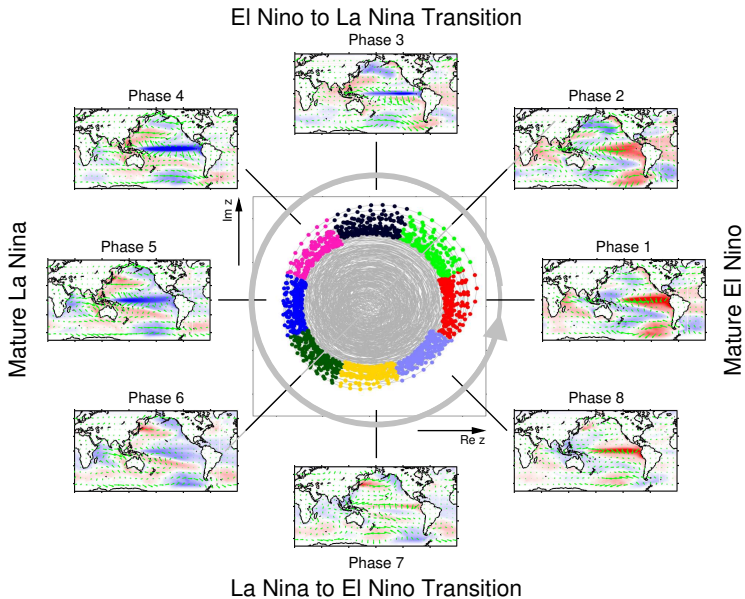


ENSO lifecycle recovered from approximate Koopman eigenfunctions

ERSSTv4 ENSO lifecycle

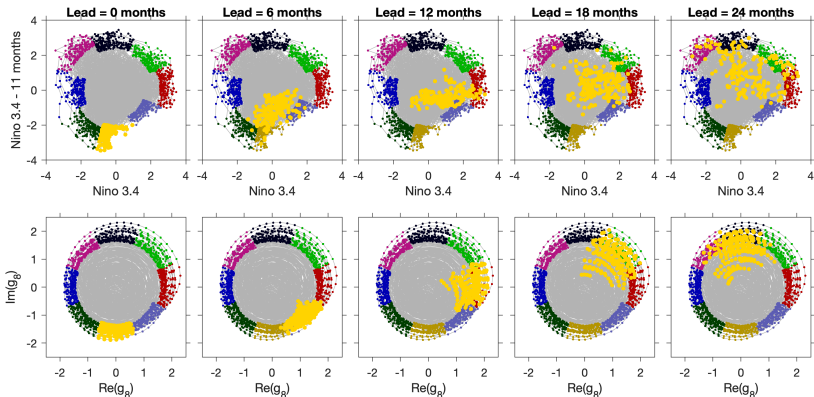


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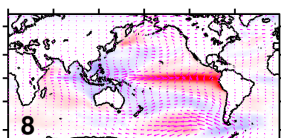
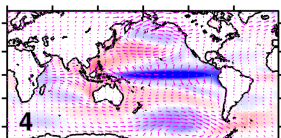
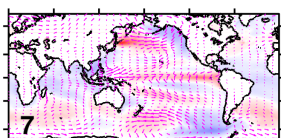
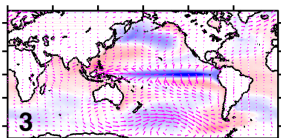
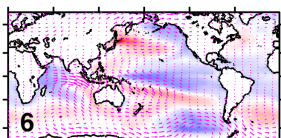
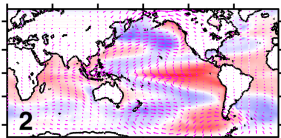
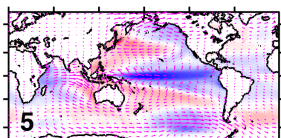
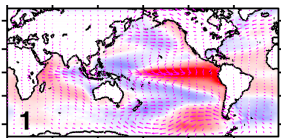


ENSO lifecycle recovered from approximate Koopman eigenfunctions

(a) CCSM4 phase evolution; Start phase = 7



ENSO phase composites



Conclusions

- **Operator-theoretic approaches**, combined with **kernel methods for machine learning**, provide a useful route to identifying coherent observables of dynamical systems.
- Eigenfunctions of kernel integral operators constructed from **delay-coordinate-mapped data** identify ϵ -approximate eigenfunctions of Koopman operators with persistent cyclical behavior, potentially under mixing (chaotic) dynamics.
- Methods are **refinable**, in the sense of spectral convergence, as the amount of training data increases.
- V&V and UQ applications, including model intercomparisons and estimation of conditional statistics, are promising directions for future work.

References

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